

Shedding Light on the Nature of Dark Matter July 13 – 24 2009

Numerical Simulations – Now and in the future

Michael Kuhlen Institute for Advanced Study Princeton, NJ

Caltech, Jul. 14th 2009

Now is the time ...

Indirect Detection

LHC



LHC might produce WIMPs!

Direct Detection

Solid Scintillation DAMA/LIBRA, KIMS

Liquid/Gas Xenon XENON-10, ZEPLIN-II Tonne-scale experiments planned

Cryogenic Detectors CDMS-II, CRESST



<u>y-rays</u> Fermi Space Telescope



A.C.T.'s (VERITAS, H.E.S.S., etc.)

anti-matter PAMELA satellite



balloon-borne: ATIC, PPB-BETS

neutrinos AMANDA, ANTARES, IceCube



1 Storenter 2 Skinner

Astronomical Surveys and

Observations

SDSS, LSST, Pan-STARRS, Keck, Gaia, SIM, TMT<u>, etc.</u>







Dark Matter Numerical Simulations

O(>10°) Particles

- > Via Lactea Project
- > Aquarius Project
- > Millennium-II
- Bolshoi



Outline of Talk

- Current state-of-the-art
 - ... of cosmological DM-only N-body simulations at the Galactic scale
- Examples of uses for DM question (past, present, and future)
 - Local Group dwarf galaxies (abundance, properties, formation history...)
 - Strong gravitational lensing effects
 - Direct detection experiments
 - Indirect detection efforts
 - Sommerfeld-enhanced subhalo annihilation signal
- Current limitations and future directions

The State-of-the-art in Cosmological DM-only N-Body Simulations





Multi-mass initial conditions



Via Lactea II – the inner 100 kpc

Whereas previous simulations were almost completely smooth in the central region, with VL-II we resolve lots of subhalos and tidal streams even down to 8 kpc.



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Via Lactea II – the inner 100 kpc

Whereas previous simulations were almost completely smooth in the central region, with VL-II we resolve lots of subhalos and tidal streams even down to 8 kpc.



The heating of stellar tidal streams holds great potential to constrain the abundance of small scale DM structure.

120

etream

V. Belokurov, IoA Cambridge / Sloan Digital Sky Survey

140

right ascension [degrees]

decli

220

200

Dwarf Spheroidal Satellite Galaxies



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| Name | D | $M_{0.3}$ | $V_{ m max}$ | $r_{ m Vmax}$ | $V_{\rm peak}$ | $r_{ m Vpeak}$ |
|-------------------|-------|---------------------------------|------------------------------|--|-----------------------|---|
| | [kpc] | $[10^7 M_{\odot}]$ | $[\rm kms^{-1}]$ | [kpc] | $[{\rm kms^{-1}}]$ | [kpc] |
| Segue 1 | 23 | $1.58^{+3.30}_{-1.11}$ | $10 \binom{17}{8.4}$ | $0.43 \begin{pmatrix} 0.89 \\ 0.29 \end{pmatrix}$ | $26\binom{55}{13}$ | $2.4 \binom{33}{1.4}$ |
| Ursa Major II | 32 | $1.09\substack{+0.89\\-0.44}$ | $13\binom{17}{11}$ | $0.59 \left({}^{0.89}_{0.31} ight)$ | $27 \binom{33}{17}$ | $3.3 \begin{pmatrix} 14\\ 2.4 \end{pmatrix}$ |
| Wilman 1 | 38 | $0.77^{+0.89}_{-0.42}$ | $8.3\left(^{11}_{7.5} ight)$ | $0.38 \left(\begin{smallmatrix} 0.62 \\ 0.29 \end{smallmatrix} \right)$ | $15\binom{27}{10}$ | $2.0 \left(\begin{smallmatrix} 3.9 \\ 0.90 \end{smallmatrix} \right)$ |
| Coma Berenices | 44 | $0.72^{+0.36}_{-0.28}$ | $9.1 \binom{12}{8.2}$ | $0.42 \begin{pmatrix} 0.62\\ 0.31 \end{pmatrix}$ | $15\binom{25}{11}$ | $1.9 \left(\begin{smallmatrix} 3.4 \\ 0.97 \end{smallmatrix} \right)$ |
| Ursa Minor | 66 | $1.79_{-0.59}^{+0.37}$ | $18\binom{21}{15}$ | $0.81 \left({{1.8}\atop{0.61}} ight)$ | $30\binom{56}{21}$ | $3.8 \binom{9.7}{2.8}$ |
| Draco | 80 | $1.87^{+0.20}_{-0.29}$ | $19\binom{22}{17}$ | $0.86 \left(\begin{smallmatrix} 2.4 \\ 0.81 \end{smallmatrix} \right)$ | $28 \binom{37}{26}$ | $3.8 \binom{32}{2.4}$ |
| Sculptor | 80 | $1.20^{+0.11}_{-0.37}$ | $13\binom{15}{12}$ | $0.64 \left(\begin{smallmatrix} 1.0 \\ 0.54 \end{smallmatrix} \right)$ | $20 \binom{25}{16}$ | $2.9 \binom{5.6}{1.6}$ |
| Sextans | 86 | $0.57\substack{+0.45 \\ -0.14}$ | $9.7 \binom{12}{8.5}$ | $0.52 \begin{pmatrix} 0.89\\ 0.37 \end{pmatrix}$ | $14\binom{19}{11}$ | $1.6 \left(\begin{smallmatrix} 3.0 \\ 0.97 \end{smallmatrix} \right)$ |
| Carina | 101 | $1.57\substack{+0.19 \\ -0.10}$ | $17\binom{22}{16}$ | $1.00 \left(\begin{smallmatrix} 2.3 \\ 0.69 \end{smallmatrix} \right)$ | $30\binom{42}{24}$ | $3.8 \binom{32}{3.3}$ |
| Ursa Major I | 106 | $1.10\substack{+0.70 \\ -0.29}$ | $14\binom{17}{13}$ | $0.84 \left(\begin{smallmatrix} 1.3 \\ 0.61 \end{smallmatrix} \right)$ | $20 \binom{30}{16}$ | $3.2 \begin{pmatrix} 6.8\\ 1.6 \end{pmatrix}$ |
| Fornax | 138 | $1.14_{-0.12}^{+0.09}$ | $15\binom{16}{14}$ | $1.1 \ \binom{1.3}{0.64}$ | $20\binom{24}{18}$ | $3.0 \binom{6.1}{1.9}$ |
| Hercules | 138 | $0.72^{+0.51}_{-0.21}$ | $11\binom{14}{9.4}$ | $0.69 \left(\begin{smallmatrix} 1.1 \\ 0.45 \end{smallmatrix} \right)$ | $14\binom{20}{12}$ | $1.9 \binom{3.8}{1.2}$ |
| Canes Venatici II | 151 | $0.70^{+0.53}_{-0.25}$ | $11\binom{13}{8.9}$ | $0.67 \left(\begin{smallmatrix} 1.1 \\ 0.44 \end{smallmatrix} ight)$ | $14 \binom{19}{11}$ | $1.8 \binom{3.7}{1.1}$ |
| Leo IV | 158 | $0.39\substack{+0.50\\-0.29}$ | $5.0 \binom{7.2}{4.2}$ | $0.35 \begin{pmatrix} 0.57 \\ 0.22 \end{pmatrix}$ | $6.7 \binom{10}{5.0}$ | $0.84 \left(\begin{smallmatrix} 1.7 \\ 0.48 \end{smallmatrix} \right)$ |
| Leo II | 205 | $1.43_{-0.15}^{+0.23}$ | $18\binom{21}{16}$ | $1.5 \begin{pmatrix} 2.1 \\ 0.93 \end{pmatrix}$ | $24 \binom{28}{19}$ | $4.1 \binom{8.2}{2.4}$ |
| Canes Venatici I | 224 | $1.40\substack{+0.18\\-0.19}$ | $18\binom{20}{16}$ | $1.5 \binom{2.1}{1.0}$ | $22 \binom{29}{18}$ | $2.9 \binom{6.1}{2.1}$ |
| Leo I | 250 | $1.45^{+0.27}_{-0.20}$ | $19\binom{21}{17}$ | $1.7 \binom{3.1}{1.1}$ | $25\binom{27}{19}$ | $2.9 \binom{6.3}{2.1}$ |
| Leo T | 417 | $1.30^{+0.88}_{-0.42}$ | $16\binom{21}{13}$ | $1.2 \left(\begin{smallmatrix} 2.4 \\ 0.85 \end{smallmatrix} \right)$ | $19\binom{26}{17}$ | $2.4 \binom{6.1}{1.6}$ |

Central masses and densities within ~300 pc are well constrained observationally and in the numerical simulations.

This will allow direct tests of CDM theory!



Dwarf Spheroidal Satellite Galaxies



| Name | D | $M_{0.3}$ | $V_{\rm max}$ | $r_{ m Vmax}$ | V_{peak} | $r_{ m Vpeak}$ |
|----------------|-------|---------------------------------|-------------------------------|--|-----------------------|--|
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Central masses and densities within ~300 pc are well constrained observationally and in the numerical simulations.

This will allow **direct tests of CDM** theory!



Madau et al. 2008

More and better measurements of stellar velocities (including proper motions?!?) in the centers of dSph's will be useful. Higher resolution simulations, too.

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The Missing Satellites "Problem"...



There appears to be a discrepancy between the number of luminous satellites and the number of dark matter halos predicted from simulation.

Solutions: Observational Completeness? Astrophysics? Particle physics?



The Subhalo Population – Tidal Mass Loss

- Subhalos orbit through host halo and are subject to tidal interactions.
- > Strongest during peri-center passage.
- > Tidal mass loss from outside in.
- > Diverse amount of tidal mass loss.

Reduces M_t, V_{max}, R_{Vmax}.





The Subhalo Population – Tidal Mass Loss



- Tidal mass loss is stronger for more massive halos (higher V_{max} @ z=1).
- > Halos with V_{max}=10 km/s retain about 40% of their mass from z=1 to today.
- > 97% of all z=1 subhalos still have an identifiable remnant at z=0.



- Subhalos are more concentrated in the inner regions.
- This due to both tidal stripping and an earlier formation time.
- > c(r=8kpc) \approx 3 \times c(field)

The Subhalo Population – Spatial Distribution



The subhalo radial distribution is anti-biased with respect to the DM density: fewer subhalos in the center.

(cf. Ghigna et al. 2000; de Lucia et al. 2004)

Depends on selection:

- strongest for M(z=0)-selected,
- weaker for Vmax(z=0)-selected,
- disappears down to ~30 kpc for peak(Vmax)-selected.

(cf. Nagai & Kravtsov 2005; Faltenbacher & Diemand 2006)

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&

The radial dependence of the number, mass, and luminosity of subhalos differs. This needs to be better understood.

radius [kpc]

Strong Gravitational Lensing Effects

(mag)

Ε





Flux ratio anomalies

(Metcalf & Madau 2001, Dalal & Kochanek 2002, etc.)



Fig. 2.—The most spectacular example of a suppressed saddle point, SDSS J0924+0219 (Inada et al. 2003). The D image should be comparable in brightness to the A image, but is actually an order of magnitude dimmer. The A and B images are minima, while C and D are saddle points. The contours are spaced by factors of 2 from the peak of the A image. The lens galaxy is seen at the center. In this infrared image, the suppression of the saddle point could be due to either microlensing or substructure.



Strong Gravitational Lensing Effects





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Time delay perturbations KEETON & MOUSTAKAS



Strong Gravitational Lensing Effects





Flux ratio anomalies

(Metcalf & Madau 2001, Dalal & Kochanek 2002, etc.)



Time delay perturbations

KEETON & MOUSTAKAS



The effects of substructures on strong lenses (flux ratio anomalies, time delay perturbations) should also be investigated with realistic DM simulations.

J0924+0219 (Inada et al. 2003). The D image should be comparable in brightness to the A image, but is actually an order of magnitude dimmer. The A and B images are minima, while C and D are saddle points. The contours are spaced by factors of 2 from the peak of the A image. The lens galaxy is seen at the center. In this infrared image, the suppression of the saddle point could be due to either microlensing or substructure.



Direct Detection

The scattering event rate (events/recoil energy) is given by:

$$\frac{dR}{dE_R} = N_T M_N \frac{\rho_{\chi}}{2m_{\chi}\mu_{ne}^2} \frac{(f_p Z + f_n (A - Z))^2}{f_n^2} F^2[E_R] \int_{\beta_{min}}^{\infty} \frac{f(v)}{v} dv$$

This depends on the local DM density ρ_{χ} and the velocity distribution function f(v).

A typical assumption is ρ_{χ} =0.3 GeV/cm³ and f(v) a Maxwellian with σ = 220 km/s truncated at an escape speed of 500-600 km/s.

Kamionkowski & Koushiappas (2008)





Vogelsberger et al. (2009) Aquarius Simulations





The DAMA controversy

An **annual modulation** of the scattering rate is expected, owing to the motion of the Earth around the Sun, which itself is moving through a Galactic DM "wind".



The DAMA collaboration has a reported an **8-\sigma detection** of this modulation and interpreted it as WIMP scattering.



Conventional DM models that fit the DAMA result appear to excluded by other experiments!



Velocity Distribution in Via Lactea/GHALO



Inelastic Dark Matter

Inelastic Dark Matter (Tucker-Smith & Weiner 2001) provides one way to **reconcile DAMA with other experiments**.

iDM is a simple extension of standard WIMP DM, with 2 new properties:

1) the DM particle χ has an excited state χ^* with m_{χ^*} - $m_{\chi} = \delta \sim 100 \text{ keV}$ 2) only inelastic scattering is permitted: $\chi \text{ N} \rightarrow \chi^* \text{ N}$

Only encounters with enough kinetic energy to excite the DM particle will lead to scattering:

$$\beta_{\min} = \sqrt{\frac{1}{2m_N E_R}} \left(\frac{m_N E_R}{\mu} + \delta\right)$$

Consequences:

- > Easier to scatter off heavier targets: favors DAMA (I, A=127) over CDMS (Ge, A=73).
- > Must sample the high-v tail of f(v), increases annual modulation: favors DAMA.
- The recoil spectrum is altered, with fewer (or no) low energy events: favors DAMA over XENON

Annual Modulation and Peak Day



Inelastic DM and Via Lactea



The shape, location, and extent of the iDM parameter space preferred by the DAMA measurement depends quite sensitively on f(v).

Global epartures from M-B and local variations can make the DAMA measurement **compatible with all current experimental limits**.

Inelastic DM and Via Lactea



Direct detection event rates can depend significantly on the local phase-space distribution of DM, especially for inelastic DM. Directionally sensitive experiments are being designed...

Indirect Detection



<u>Neutrinos</u>



Positrons

PAMELA





Indirect Detection

DM (WIMP) annihilation signal





3 ways to get gamma-rays:





Angle for 68% containment (degrees) 10000 andoff thin section (best osf) 10 violif thin section (best osf handoff thick section handoff (thick+thin) +++++ handoff thick section 8000 Area (cm^2) 1Ē 10 2000 10 10² 103 104 10² 10³ 10⁴ Energy (MeV) 10 Energy (MeV)

Detector properties

105

$$N_{\gamma} = \left[\int_{\text{line of sight}}^{\rho_{\text{DM}}^{2}} \frac{dl(\psi)}{2M_{\chi}^{2}} \right] \frac{\langle \sigma v \rangle}{2M_{\chi}^{2}} \left[\int_{E_{th}}^{M_{\chi}} \left(\frac{dN_{\gamma}}{dE} \right)_{\text{SUSY}}^{A_{\text{eff}}}(E) dE \right] \frac{\Delta \Omega}{4\pi} \tau_{\text{exp}}$$



Simulated Dark Matter Annihilation Map



Galactic Center: GR Point Sources



Backgrounds: Galactic GR Background



Backgrounds: Extragalactic GR Background



The Total Signal



Known MW dSph Satellites as Annihilation Sources



Known MW dSph Satellites as Annihilation Sources



Known MW dSph Satellites as Annihilation Sources





| Fornax | 138 | $3.5 \begin{pmatrix} 4.4 \\ 3.0 \end{pmatrix}$ | $2.9 \begin{pmatrix} 3.3\\ 2.3 \end{pmatrix}$ | $1.4 \begin{pmatrix} 1.8\\ 1.3 \end{pmatrix}$ | $1.00 \left(\begin{smallmatrix} 1.2 \\ 0.74 \end{smallmatrix} \right)$ |
|------------------|-----|--|--|---|--|
| Sextans | 86 | $1.2 \left(\begin{smallmatrix} 2.0 \\ 0.77 \end{smallmatrix} \right)$ | $1.1 \begin{pmatrix} 1.8\\ 0.69 \end{pmatrix}$ | $1.3 \left(\begin{smallmatrix} 2.1 \\ 0.83 \end{smallmatrix} \right)$ | $0.86 \binom{1.4}{0.55}$ |
| Leo II | 205 | $4.6 \begin{pmatrix} 6.5\\ 3.8 \end{pmatrix}$ | $3.1 \begin{pmatrix} 4.7\\ 2.1 \end{pmatrix}$ | $0.88 \left(\begin{smallmatrix} 1.2 \\ 0.73 \end{smallmatrix} \right)$ | $0.55 \begin{pmatrix} 0.85\\ 0.37 \end{pmatrix}$ |
| Canes Venatici I | 224 | $4.6 \begin{pmatrix} 7.9\\ 3.8 \end{pmatrix}$ | $3.1 \binom{5.0}{2.3}$ | $0.73 \left(\begin{smallmatrix} 1.3 \\ 0.60 \end{smallmatrix} \right)$ | $0.48 \left(\begin{smallmatrix} 0.79 \\ 0.35 \end{smallmatrix} \right)$ |
| | | .7.0. | . # 4. | .1.0 . | .0.72. |

Indirect Detection: the leptonic signal

There are currently at least four unexplained observations of energetic electrons and positrons in the Milky Way halo.

1) The WMAP "haze"

3) The PAMELA measurement of a positron excess at E > 10 GeV.

- 2) The INTEGRAL/SPI 511 keV Galactic Center signal
- 4) The ATIC and PPB-BETS measurements of an (e⁺+e⁻) bump at E ~ 500 – 800 GeV.



Indirect Detection: the leptonic signal

There are currently at least four unexplained observations of energetic electrons and positrons in the Milky Way halo.

1) The WMAP "haze"

3) The PAMELA measurement of a positron excess at E > 10 GeV.

- 2) The INTEGRAL/SPI 511 keV Galactic Center signal
- 4) Fermi does not see the ATIC/PPB-BETS bump. Lack of feature implies $M_{\chi} > \sim 1$ TeV.



Could this be DM annihilation?

A typical $M_{\chi} \sim 1$ TeV WIMP DM particle would need to have a cross section ~ 1000 larger than the thermal relic value to explain the PAMELA and ATIC/PPB_BETS/Fermi results.

The local DM distribution is **not clumpy enough** to give such a large boost factor.



Extrapolating the VL-II subhalo abundance in mass, taking into account radial anti-bias, gives a local boost factor of B = 1.4.

In VL-II about 1% of locations at 8kpc happen to be close enough to a subhalo to give B~10.

Need some other mechanism to increase the cross section → Sommerfeld enhancement?

Could this be DM annihilation?

A typical $M_{\chi} \sim 1$ TeV WIMP DM particle would need to have a huge cross section to explain the PAMELA and ATIC/PPB_BETS/Fermi results.

With Sommerfeld enhancement and a ~TeV particle, the data can be nicely fit:



This seems to me to be a stalemate. Have local anti-particle flux measurements told us all they can?





But so can, for example, pulsars...

Example: Sommerfeld-enhanced subhalo annihilation signal



Three regimes: (cf. Lattanzi & Silk 2008) I. Large velocities (ϵ_v >1): no enhancement II. Intermediate velocities ($\epsilon_{\phi} < \epsilon_v < 1$): $S \sim \pi \frac{\alpha}{v}$ III. Low velocities ($\epsilon_v < \epsilon_{\phi}$), depends on m_x: either saturation: S = $1/\epsilon_{\phi}$ or resonances (bound states) at some values of m_x: $S \sim \alpha \frac{m_{\phi}}{m_{\chi}} \frac{1}{v^2}$



Velocities in Via Lactea II



Velocity distribution in host and subhalos and resulting Sommerfeld enhancement factors



Velocities in Via Lactea II



Velocity distribution in host and subhalos and resulting Sommerfeld enhancement factors



With Sommerfeld Enhancement

Model A

Model A

Three regimes: (cf. Lattanzi & Silk 2008) I. Large velocities ($\epsilon_v > 1$): no enhancement II. Intermediate velocities ($\epsilon_{\phi} < \epsilon_v < 1$): $S \sim \pi \frac{\alpha}{v}$ III. Low velocities ($\epsilon_v < \epsilon_{\phi}$), depends on m_x : either saturation: $S = 1/\epsilon_{\phi}$ or resonances (bound states) at some values of m_x : $S \sim \alpha \frac{m_{\phi}}{m_{\chi}} \frac{1}{v^2}$



With Sommerfeld Enhancement

The Sommerfeld enhancement also brightens the central regions of subhalos!



Many of these Sommerfeld models will be strongly constrained by Fermi data.

| Table 2: Detectable Subhalos. The number of subhalos that would be detected with $> 5\sigma$ significance by <i>Fermi</i> |
|---|
| after 1, 2, 5, and 10 years in orbit, for different Sommerfeld-enhanced dark matter particle models. In the two |
| right-most columns we give the median distance and mass of the detectable clumps after 5 years in orbit. |

| Model | m_{χ} | m_{ϕ} | $\alpha \times 100$ | S_{\max} | $v_{\rm sat}$ |
|--------|------------|------------|---------------------|------------|---------------------------|
| | (TeV) | (GeV) | | | $({\rm km}~{\rm s}^{-1})$ |
| LS-1 | 4.30 | 90 | 3.307 | 1,500 | 80 |
| LS-2 | 4.45 | 90 | 3.297 | 12,000 | 28 |
| LS-3 | 4.50 | 90 | 3.288 | 70,000 | 12 |
| LS-4 | 4.55 | 90 | 3.281 | 430,000 | 4.7 |
| MPV-1a | 1.0 | 0.2 | 4.000 | 3,000 | 7.4 |
| MPV-1b | 1.0 | 0.2 | 3.739 | 16,000 | 2.4 |
| MPV-2a | 0.25 | 0.2 | 4.000 | 480 | 40 |
| MPV-2b | 0.25 | 0.2 | 4.500 | 40,000 | 3.3 |

| Model | 1 yr | 2 yr | 5 yr | 10 yr | \tilde{D} | $\tilde{M}_{ m sub}$ |
|--------|-------|------|------|-------|-------------|----------------------|
| | | | | | (kpc) | (M_{\odot}) |
| LS-1 | 12 | 19 | 29 | 38 | 24 | $1.4	imes10^7$ |
| LS-2 | 72 | 99 | 167 | 244 | 42 | $9.5 	imes 10^6$ |
| LS-3 | 225 | 311 | 457 | 583 | 56 | $6.2	imes10^6$ |
| LS-4 | 410 | 528 | 730 | 919 | 66 | $4.9	imes10^6$ |
| MPV-1a | a 5 | 7 | 12 | 15 | 16 | $9.8 	imes 10^7$ |
| MPV-11 | 9 | 14 | 25 | 36 | 25 | $4.4	imes10^6$ |
| MPV-2a | a 12 | 18 | 29 | 38 | 24 | $1.4 	imes 10^7$ |
| MPV-21 | 0 187 | 254 | 397 | 518 | 55 | $4.5 	imes 10^6$ |

For DM-only simulations like the ones I discussed:

- 1) We're only beginning to **resolve the ~100 parsec scale** in subhalos
 - Better match between observationally constrained properties and simulated subhalos: M_{0.3}, central (phase-space) densities.

2) Local phase-space structure at ~8kpc has not yet converged

- Implications for direct detection (especially for inelastic DM), and local DM annihilation (positron fraction, e⁺e⁻ flux, neutrinos, ...).
- · Relevance for stellar tidal streams?

3) Cosmic variance?

- How representative is any one high resolution numerical simulation?
- 4) More analysis of the existing simulations is needed
 - $\cdot\,$ time evolution, subhalo merger trees, tidal streams, etc.

Absence of Baryonic Physics!!!



Absence of Baryonic Physics!!!

Where is it most critical and likely to change things?

1) For the Galactic Center and the smooth host halo profile

- $\cdot\,$ Gas cooling and settling drags DM to center: adiabatic contraction
- Dynamical friction of infalling satellites with dense baryonic condensations can remove DM from center
 Romano-Diaz et al. 2008
- So can stellar bars and a binary SMBH...



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 - So can stellar bars and a binary SMBH...

2) Increase substructure survivability and hence spatial distribution

- · Central baryonic condensation might make them more resilient to tidal stripping.
- · Greater retained mass means more efficient dynamical friction.

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- 1) For the Galactic Center and the smooth host halo profile
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2) Increase substructure survivability and hence spatial distribution

- · Central baryonic condensation might make them more resilient to tidal stripping.
- · Greater retained mass means more efficient dynamical friction.
- 3) Increase disk heating? (e.g. Kazantzidis et al. 2009)

Current Limitations and Romano-Diaz et al. 2008

Absence of Baryonic Physics!!!

Where is it most critical and like

- 1) For the Galactic Center and the $\overline{1}$
 - \cdot Gas cooling and settling drags DM to E
 - Dynamical friction of infalling satellites remove DM from center
 - \cdot So can stellar bars and a binary SMBH

2) Increase substructure survivabili

- Central baryonic condensation might r
- Greater retained mass means more ef
- 3) Increase disk heating? (e.g. Kazar
- 4) Modify DM velocity dispersions in the host halo, in subhalos?
 - Implications for Sommerfeld-enhanced models



Current Limitations and Romano-Diaz et al. 2008

Absence of Baryonic Physics!!!

Where is it most critical and like

- 1) For the Galactic Center and the $\hat{\mathbf{f}_{\omega}}^{10^3}$
 - Gas cooling and settling drags DM to $c_{\underline{\xi}}$
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 - $\cdot\,$ So can stellar bars and a binary SMBH
- 2) Increase substructure survivabilit
 - Central baryonic condensation might m

Baryonic physics must be included, but

- It's computationally much more expensive
- It makes the results dependent on subgrid physics, ad-hoc assumptions about gas cooling, star formation, feedback, etc.

